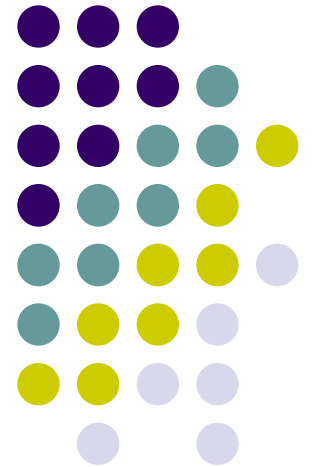
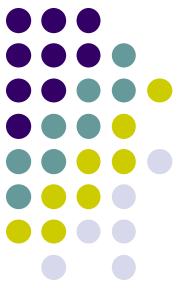


Chapter 25

Electric Potential





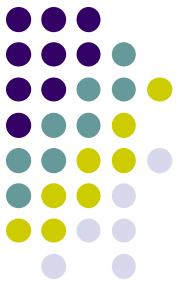
Electrical Potential Energy

- When a **test charge** is placed in **an electric field**, it experiences **a force**

$$\vec{F} = q_o \vec{E}$$

- The force is conservative
- If the **test charge** is moved in the field by **some external agent**, the **work done by the field is the negative** of the work done by the external agent
- $d\vec{s}$ is an **infinitesimal displacement vector** that is **oriented tangent** to a path through space

Electric Potential Energy, cont



- The work done by the electric field is

$$\vec{F} \cdot d\vec{s} = q_o \vec{E} \cdot d\vec{s}$$

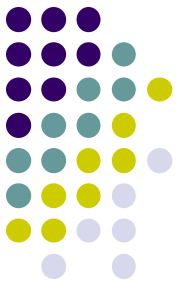
- As this work is done by the field, the potential energy of the charge-field system is changed

$$\text{by } \Delta U = -q_o \vec{E} \cdot d\vec{s}$$

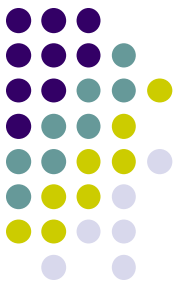
- For a finite displacement of the charge from A to B,

$$\Delta U = U_B - U_A = -q_o \int_A^B \vec{E} \cdot d\vec{s}$$

Electric Potential Energy, final



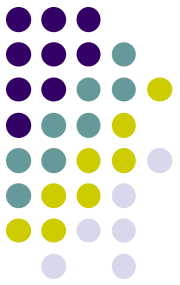
- Because the force is conservative, **the line integral does not depend on the path** taken by the charge
- This is **the change in potential energy** of the system



Electric Potential

- The potential energy per unit charge, U/q_o , is the **electric potential**
 - **The potential** is characteristic **of the field only**
 - **The potential energy** is characteristic of the **charge-field system**
 - The potential is independent of the value of q_o
 - The potential has a value at every point in an electric field
- **The electric potential is**

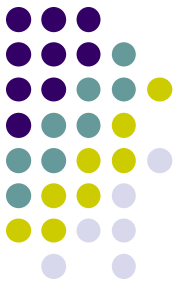
$$V = \frac{U}{q_o}$$



Electric Potential, cont.

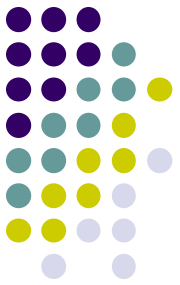
- The potential is a scalar quantity
 - Since energy is a scalar
- As a charged particle moves in an electric field, it will experience a change in potential

$$\Delta V = \frac{\Delta U}{q_o} = -\int_A^B \vec{\mathbf{E}} \cdot d\vec{\mathbf{s}}$$



Electric Potential, final

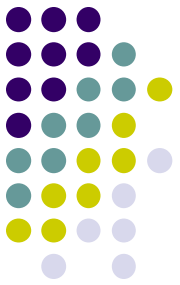
- The ***difference in potential is the meaningful quantity***
- We often take **the value of the potential to be zero at some convenient point** in the field
- **Electric potential is a scalar characteristic of an electric field,** independent of any charges that may be placed in the field



Work and Electric Potential

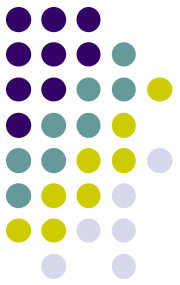
- Assume a charge moves in an electric field without any change in its kinetic energy
- The work **W** performed on the charge is

$$W = \Delta U = q \Delta V$$



Units

- $1 \text{ V} = 1 \text{ J/C}$
 - V is a volt
 - It takes one joule of work to move a 1-coulomb charge through a potential difference of 1 volt
- In addition, $1 \text{ N/C} = 1 \text{ V/m}$
 - This indicates we can interpret the electric field as a measure of the rate of change with position of the electric potential



Electron-Volts

- Another **unit of energy** that is commonly used in atomic and nuclear physics is **the electron-volt**
- **One *electron-volt*** is defined as **the energy a charge-field system gains or loses** when **a charge of magnitude e** (an electron or a proton) is **moved through a potential difference of 1 volt**
 - $1 \text{ eV} = 1.60 \times 10^{-19} \text{ J}$

Potential Difference in a Uniform Field



- The equations for electric potential can be simplified if the electric field is **uniform**:

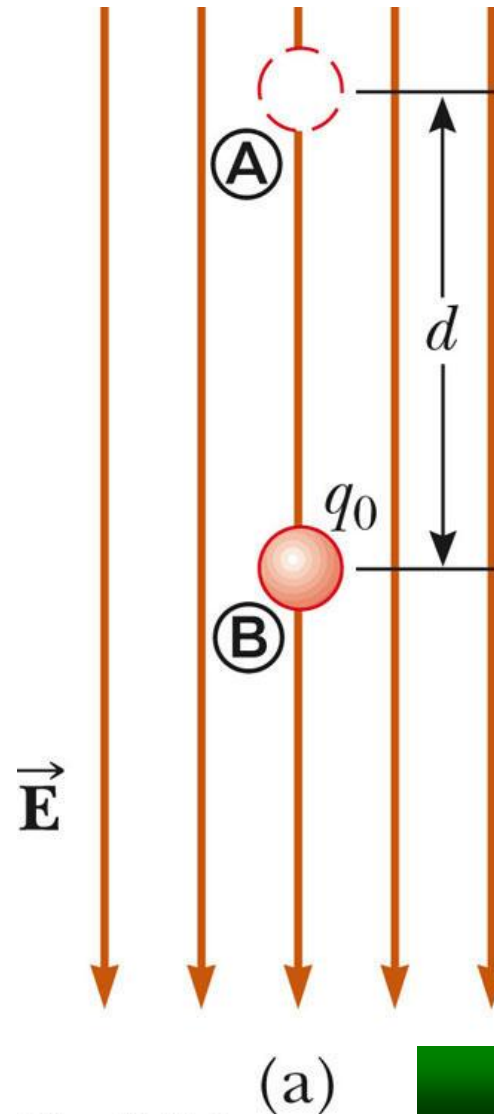
$$V_B - V_A = \Delta V = -\int_A^B \vec{\mathbf{E}} \cdot d\vec{\mathbf{s}} = -E \int_A^B ds = -Ed$$

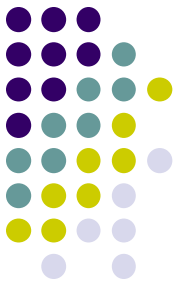
- The **negative sign** indicates that the electric potential at point B is lower than at point A
- Electric field lines always **point in the direction of decreasing electric potential**

Energy and the Direction of Electric Field



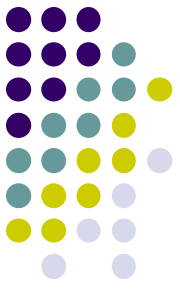
- When the **electric field is directed downward**, **point B** is at a lower potential than point **A**
- When a **positive test charge moves from A to B**, the charge-field system **loses potential energy**
- Use the active figure to compare the motion in the electric field to the motion in a gravitational field





More About Directions

- A system consisting of a positive charge and an electric field **loses** electric potential energy when the charge moves **in the direction of the field**
 - An electric field does work on a positive charge when the charge moves in the direction of the electric field
- The **charged particle gains kinetic energy** equal to the **potential energy lost by the charge-field system**
 - Another example of Conservation of Energy

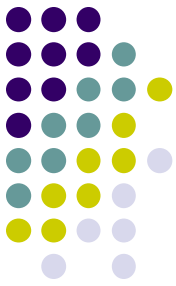
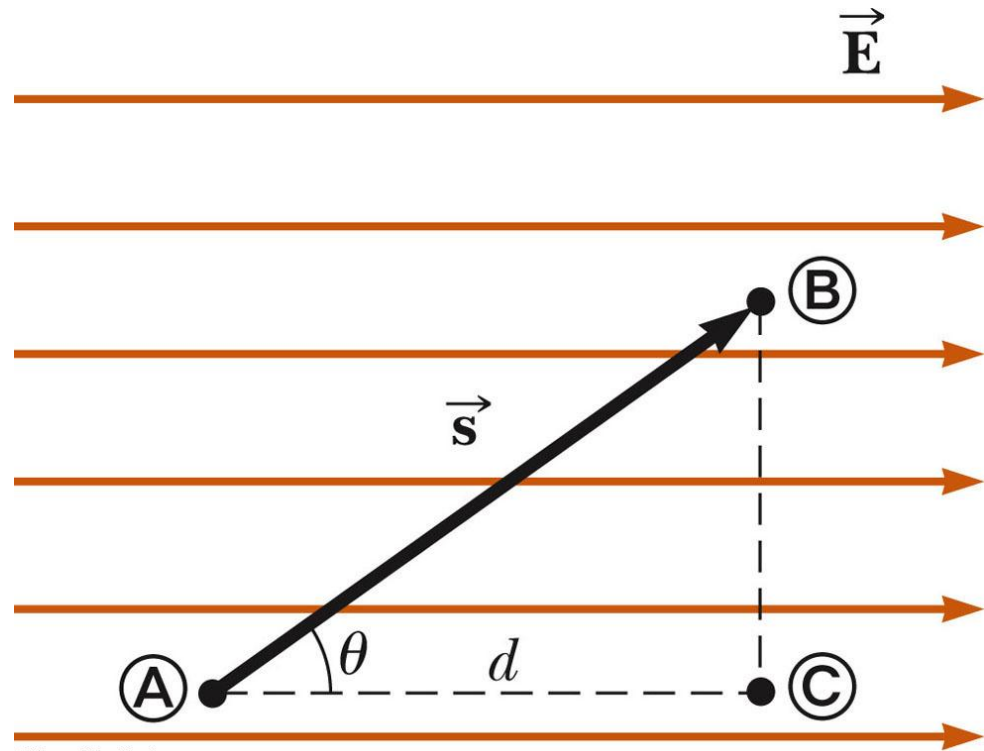


Directions, cont.

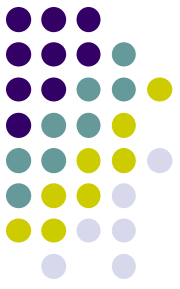
- If q_0 is negative, then ΔU is positive
- A system consisting of a negative charge and an electric field *gains* potential energy when the **charge moves in the direction of the field**
- In order for a negative charge to move in the direction of the field, an external agent must do positive work on the charge

Equipotentials

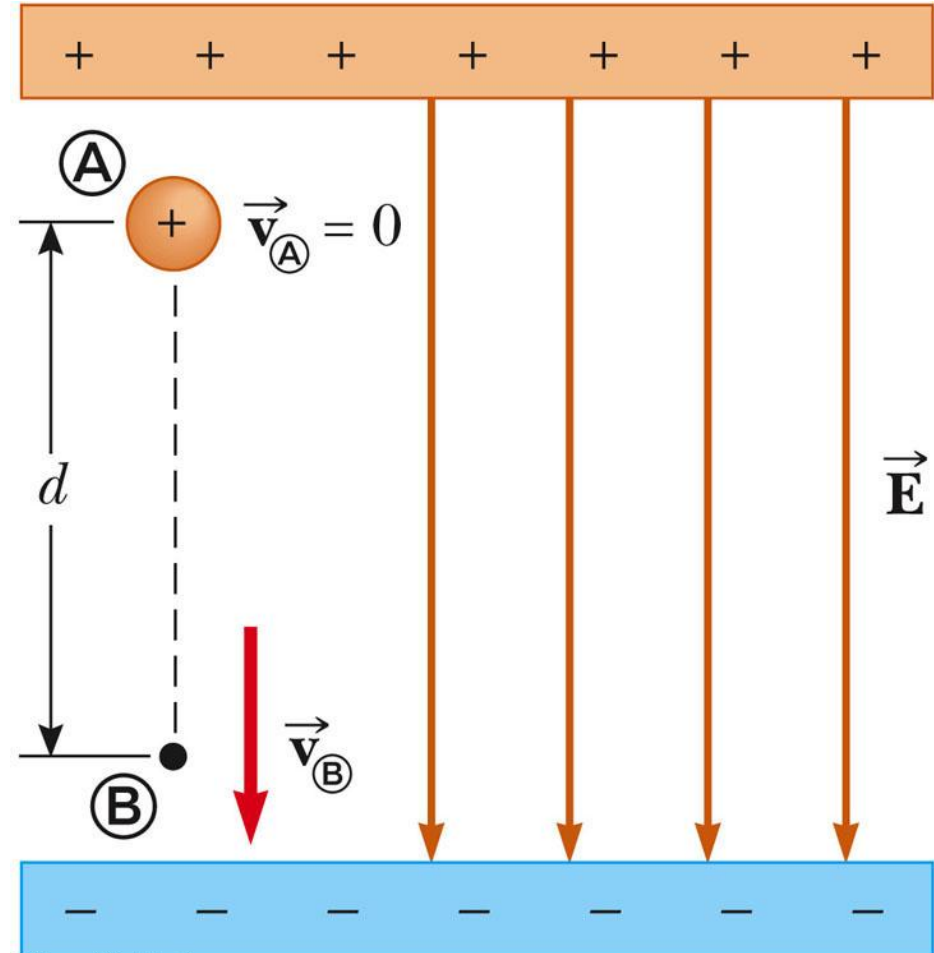
- Point B is at a lower potential than point A
- Points B and C are at the same potential
 - All points in a plane perpendicular to a uniform electric field are at the same electric potential
- The name **equipotential surface** is given to any surface consisting of a continuous distribution of points **having the same electric potential**

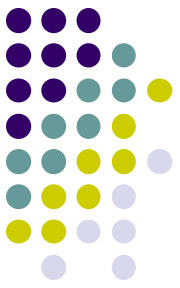


Charged Particle in a Uniform Field, Example



- A positive charge is released from rest and moves in the direction of the electric field
- The change in potential is negative
- The change in potential energy is negative
- The force and acceleration are in the direction of the field
- Conservation of Energy can be used to find its speed





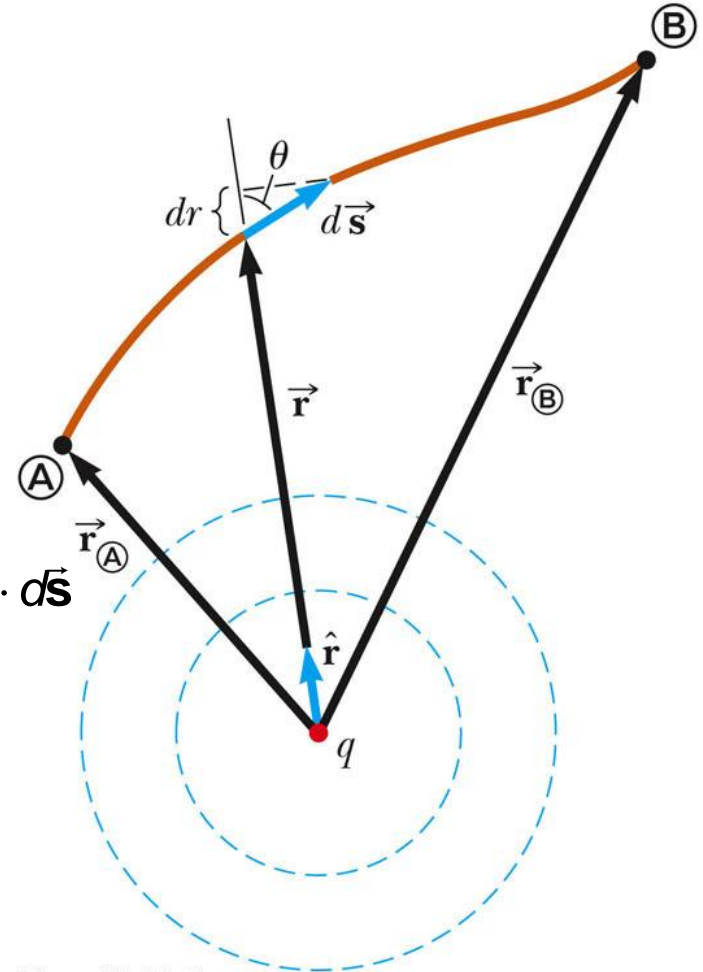
Potential and Point Charges

- A positive point charge produces a field directed radially outward
- The potential difference between points A and B will be:

$$V_B - V_A = k_e q \left[\frac{1}{r_B} - \frac{1}{r_A} \right]$$

Integrate:

$$\Delta V = \frac{\Delta U}{q_o} = - \int_A^B \vec{E} \cdot d\vec{s}$$



Potential and Point Charges, cont.



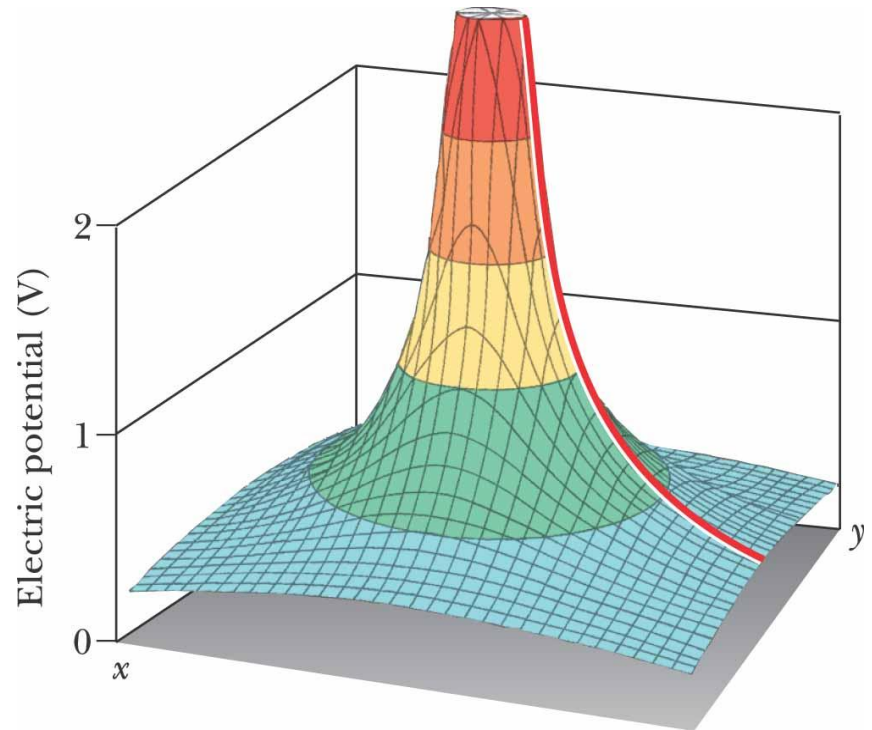
- The electric potential is independent of the path between points A and B
- It is customary to choose a reference potential of $V = 0$ at $r_A = \infty$
- Then the potential at some point r is

$$V = k_e \frac{q}{r}$$

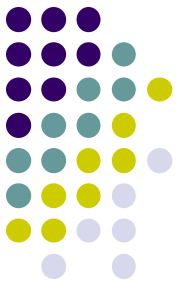
Electric Potential of a Point Charge



- The electric potential in the plane around a single point charge is shown
- The red line shows the $1/r$ nature of the potential



Electric Potential with Multiple Charges



- The electric potential due to several point charges is the sum of the potentials due to each individual charge
 - This is another example of the superposition principle
 - The sum is the algebraic sum

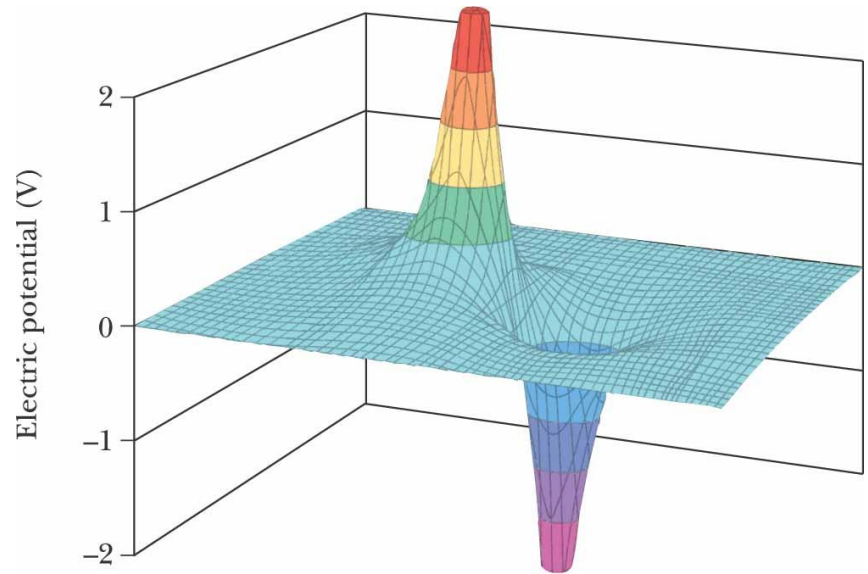
$$V = k_e \sum_i \frac{q_i}{r_i}$$

- $V = 0$ at $r = \infty$

Electric Potential of a Dipole



- The graph shows the potential (y-axis) of an electric dipole
- The steep slope between the charges represents the strong electric field in this region



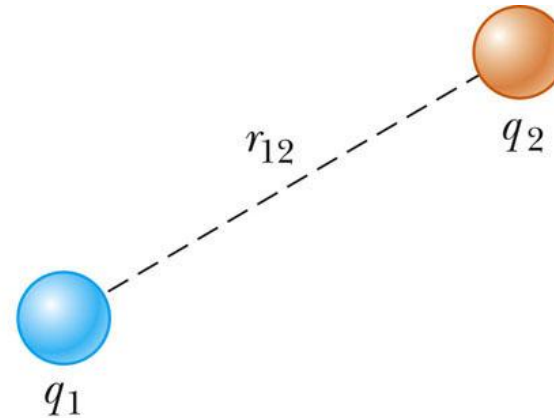
Potential Energy of Multiple Charges



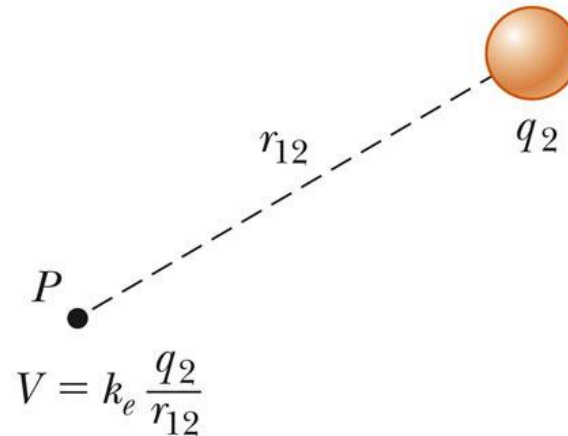
- Consider two charged particles
- The potential energy of the system is

$$U = k_e \frac{q_1 q_2}{r_{12}}$$

- Use the active figure to move the charge and see the effect on the potential energy of the system

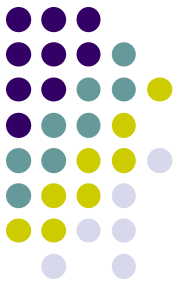


(a)

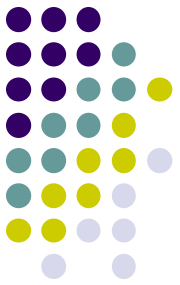


(b)

More About U of Multiple Charges



- If the two charges are the same sign, U is positive and work must be done to bring the charges together
- If the two charges have opposite signs, U is negative and work is done to keep the charges apart

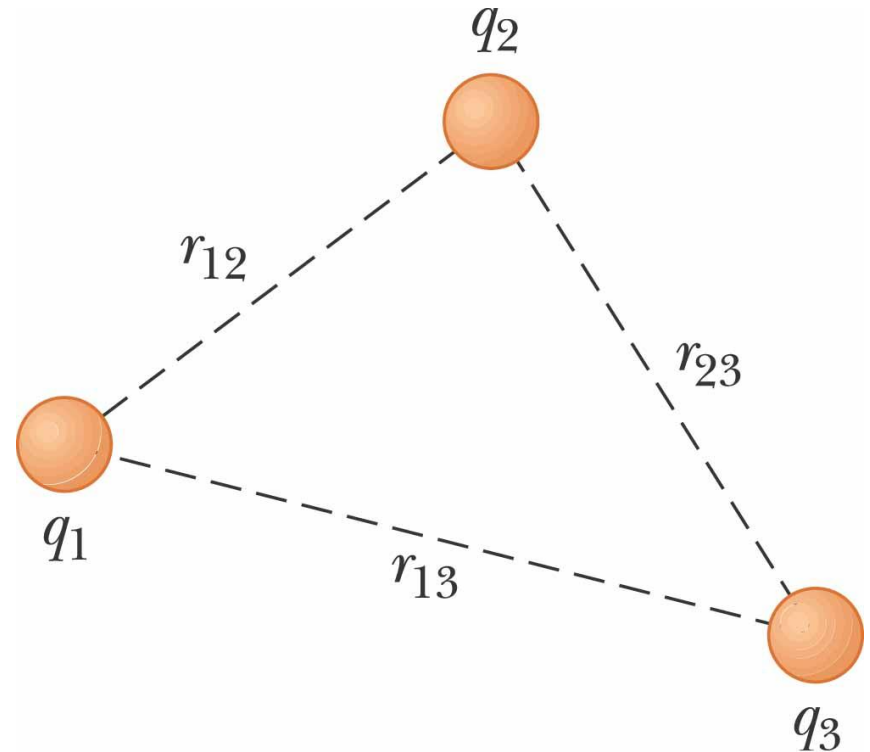


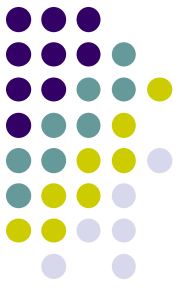
U with Multiple Charges, final

- If there are more than two charges, then find U for each pair of charges and add them
- For three charges:

$$U = k_e \left(\frac{q_1 q_2}{r_{12}} + \frac{q_1 q_3}{r_{13}} + \frac{q_2 q_3}{r_{23}} \right)$$

- The result is independent of the order of the charges





Finding E From V

- Assume, to start, that the field has only an x component

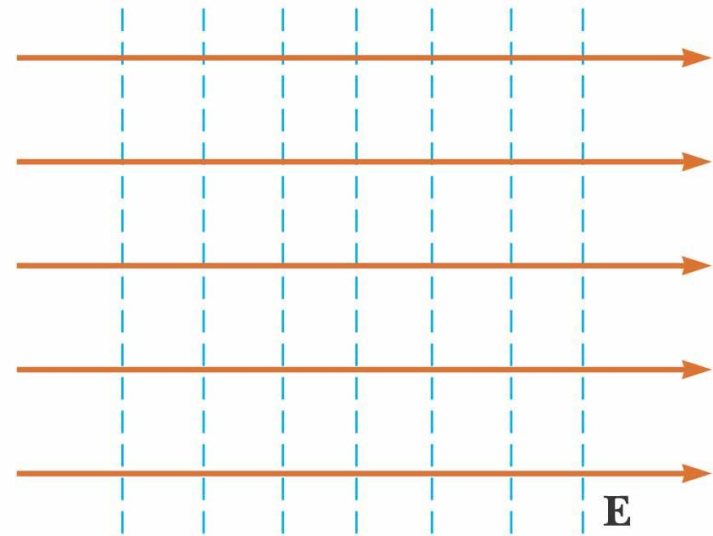
$$E_x = -\frac{dV}{dx}$$

- Similar statements would apply to the y and z components
- Equipotential surfaces must always be perpendicular to the electric field lines passing through them

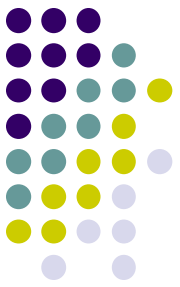
E and V for an Infinite Sheet of Charge



- The equipotential lines are the dashed blue lines
- The electric field lines are the brown lines
- The equipotential lines are everywhere perpendicular to the field lines

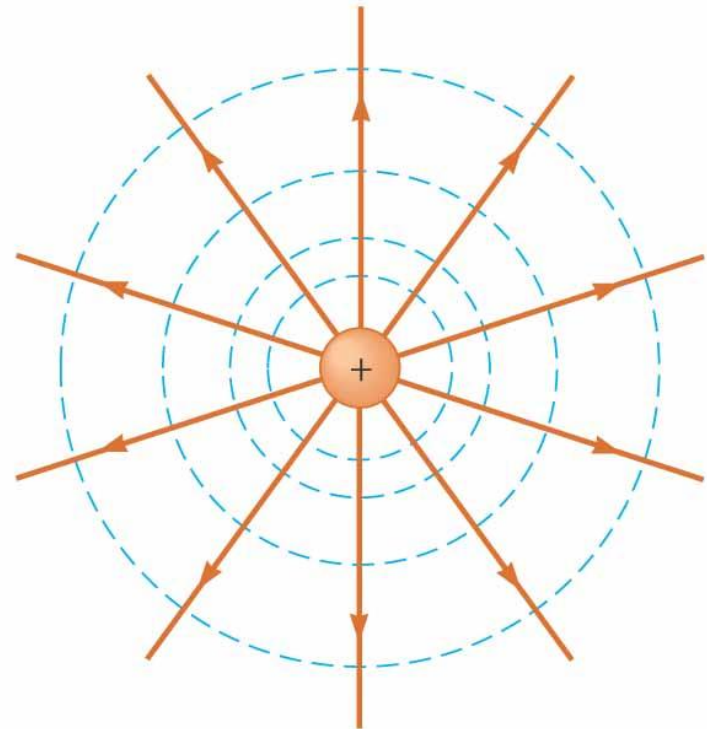


(a)



E and V for a Point Charge

- The equipotential lines are the dashed blue lines
- The electric field lines are the brown lines
- The equipotential lines are everywhere perpendicular to the field lines

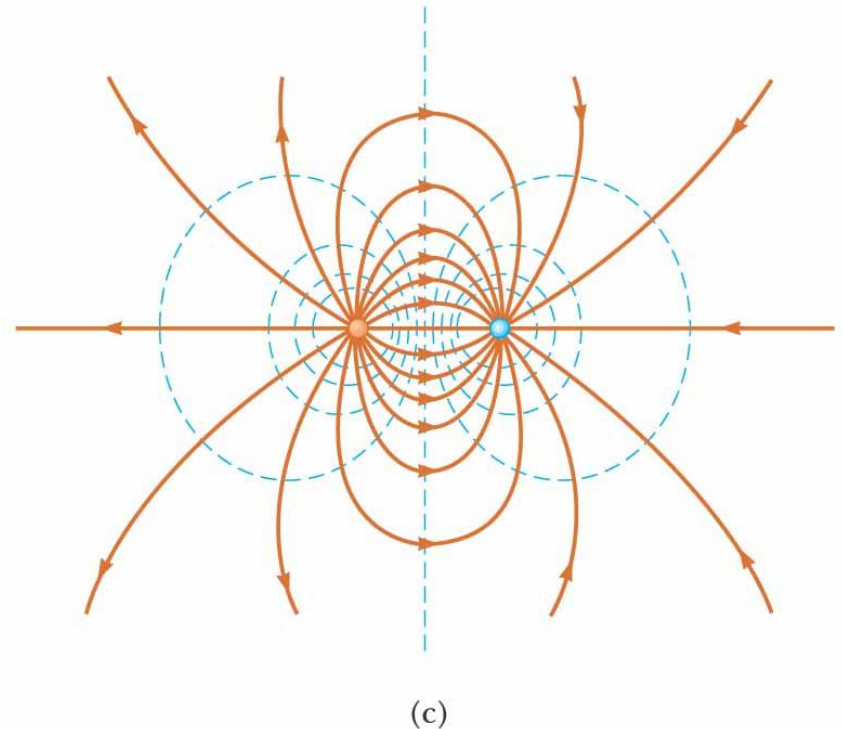


(b)

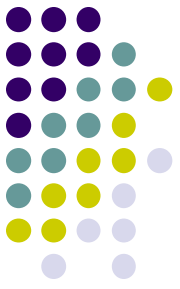
E and V for a Dipole



- The equipotential lines are the dashed blue lines
- The electric field lines are the brown lines
- The equipotential lines are everywhere perpendicular to the field lines



Electric Field from Potential, General



- In general, the electric potential is a function of all three dimensions
- Given $V(x, y, z)$ you can find E_x , E_y and E_z as partial derivatives

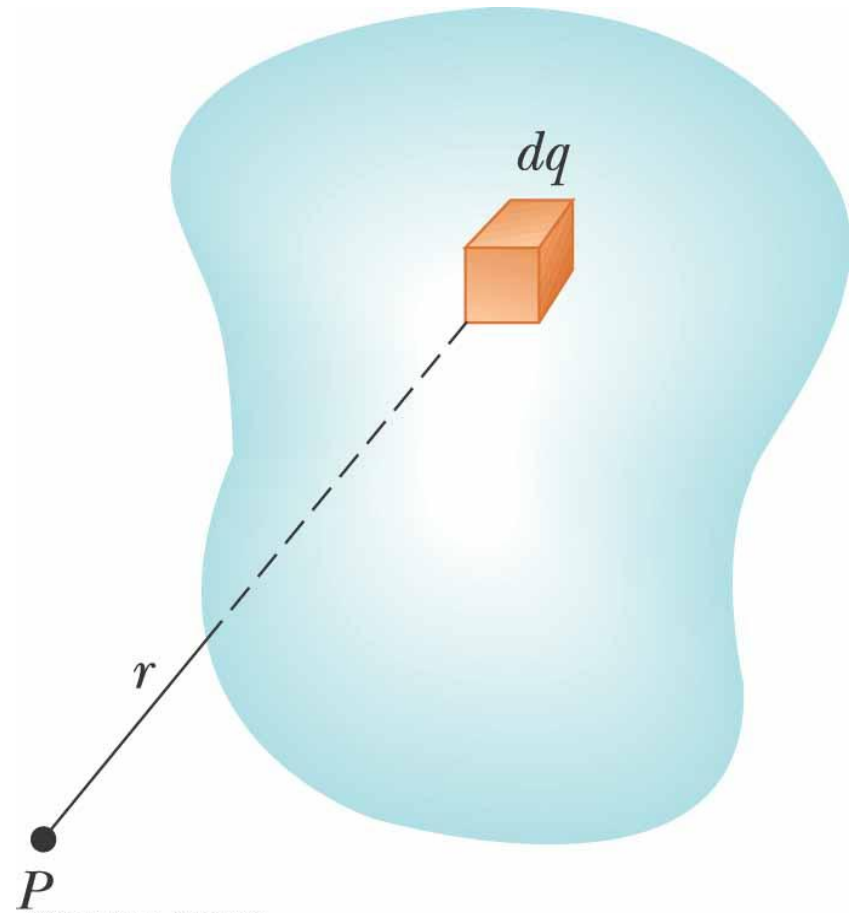
$$E_x = -\frac{\partial V}{\partial x} \quad E_y = -\frac{\partial V}{\partial y} \quad E_z = -\frac{\partial V}{\partial z}$$

Electric Potential for a Continuous Charge Distribution

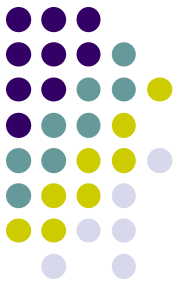


- Consider a small charge element dq
 - Treat it as a point charge
- The potential at some point due to this charge element is

$$dV = k_e \frac{dq}{r}$$



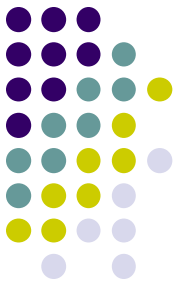
V for a Continuous Charge Distribution, cont.



- To find the total potential, you need to integrate to include the contributions from all the elements

$$V = k_e \int \frac{dq}{r}$$

- This value for V uses the reference of $V = 0$ when P is infinitely far away from the charge distributions



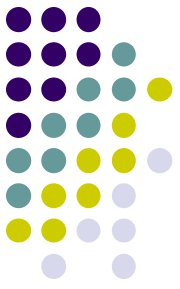
V From a Known E

- If the electric field is already known from other considerations, the potential can be calculated using the original approach

$$\Delta V = -\int_A^B \vec{\mathbf{E}} \cdot d\vec{\mathbf{s}}$$

- If the charge distribution has sufficient symmetry, first find the field from Gauss' Law and then find the potential difference between any two points
 - Choose $V = 0$ at some convenient point

Problem-Solving Strategies



- *Conceptualize*
 - Think about the individual charges or the charge distribution
 - Imagine the type of potential that would be created
 - Appeal to any symmetry in the arrangement of the charges
- *Categorize*
 - Group of individual charges or a continuous distribution?

Problem-Solving Strategies, 2



- *Analyze*
 - General
 - Scalar quantity, so no components
 - Use algebraic sum in the superposition principle
 - Only changes in electric potential are significant
 - Define $V = 0$ at a point infinitely far away from the charges
 - If the charge distribution extends to infinity, then choose some other arbitrary point as a reference point

Problem-Solving Strategies, 3



- *Analyze, cont*
 - If a group of individual charges is given
 - Use the superposition principle and the algebraic sum
 - If a continuous charge distribution is given
 - Use integrals for evaluating the total potential at some point
 - Each element of the charge distribution is treated as a point charge
 - If the electric field is given
 - Start with the definition of the electric potential
 - Find the field from Gauss' Law (or some other process) if needed

Problem-Solving Strategies, final



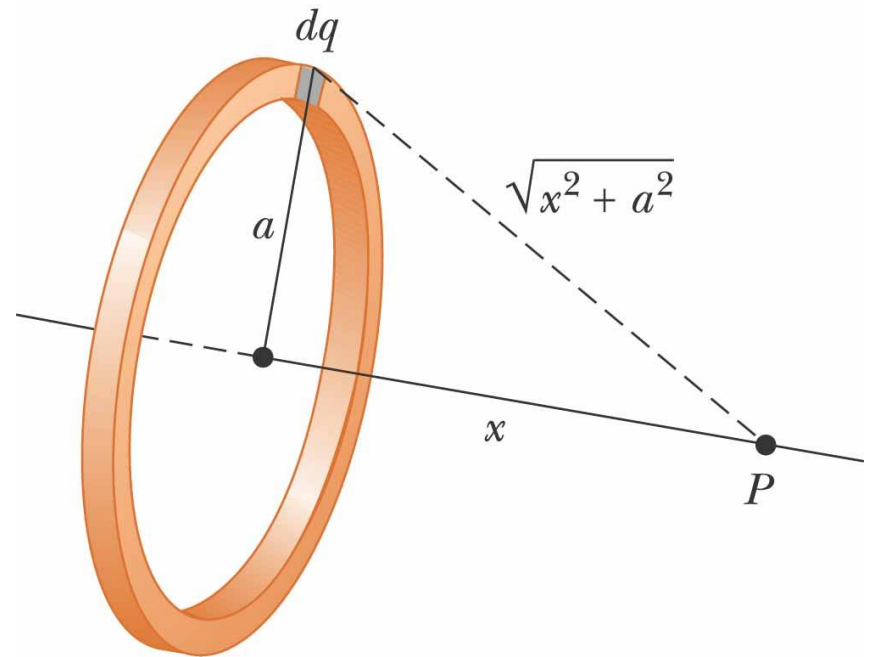
- *Finalize*
 - Check to see if the expression for the electric potential is consistent with your mental representation
 - Does the final expression reflect any symmetry?
 - Image varying parameters to see if the mathematical results change in a reasonable way

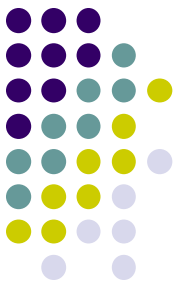
V for a Uniformly Charged Ring



- P is located on the perpendicular central axis of the uniformly charged ring
 - The ring has a radius a and a total charge Q

$$V = k_e \int \frac{dq}{r} = \frac{k_e Q}{\sqrt{a^2 + x^2}}$$

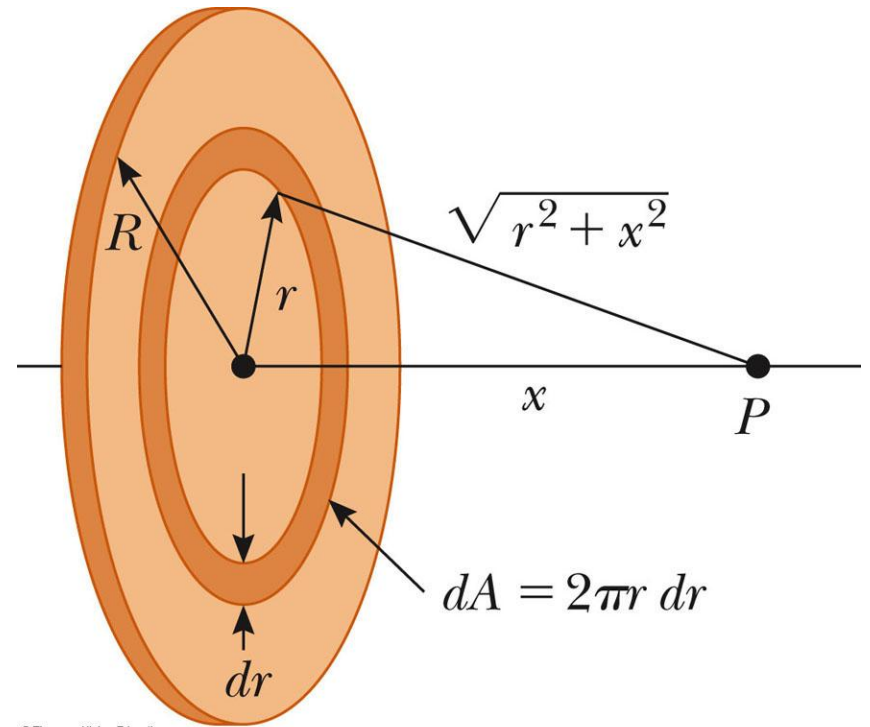


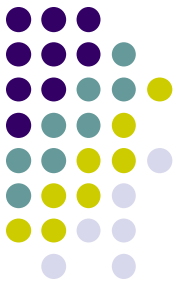


V for a Uniformly Charged Disk

- The ring has a radius R and surface charge density of σ
- P is along the perpendicular central axis of the disk

$$V = 2\pi k_e \sigma \left[\left(R^2 + x^2 \right)^{1/2} - x \right]$$

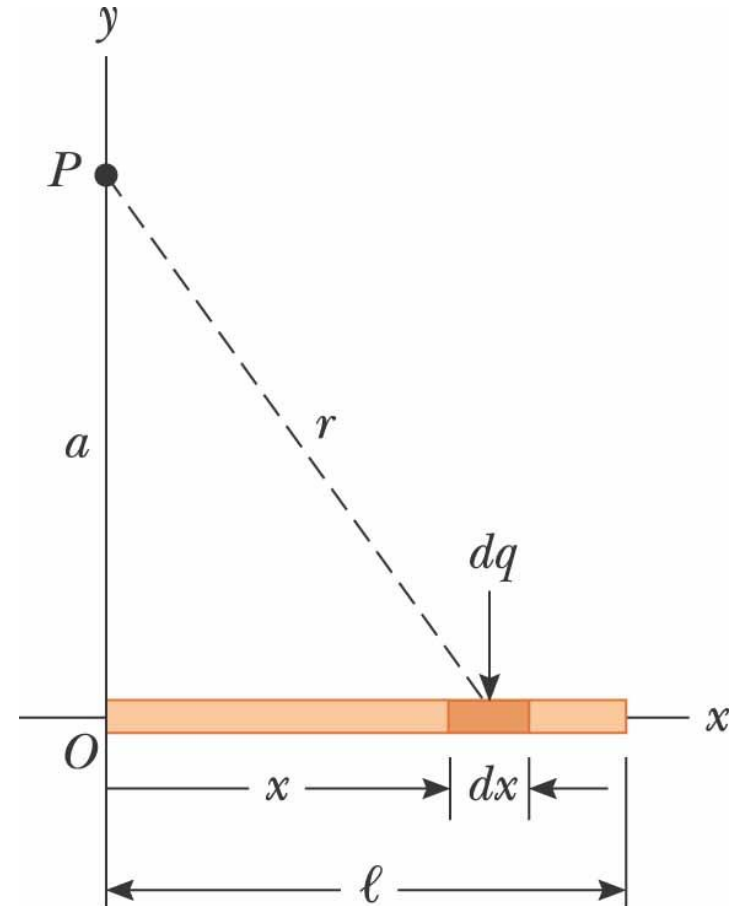


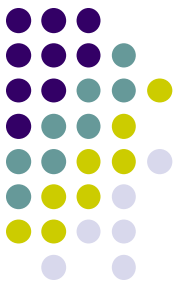


V for a Finite Line of Charge

- A rod of length ℓ has a total charge of Q and a linear charge density of λ

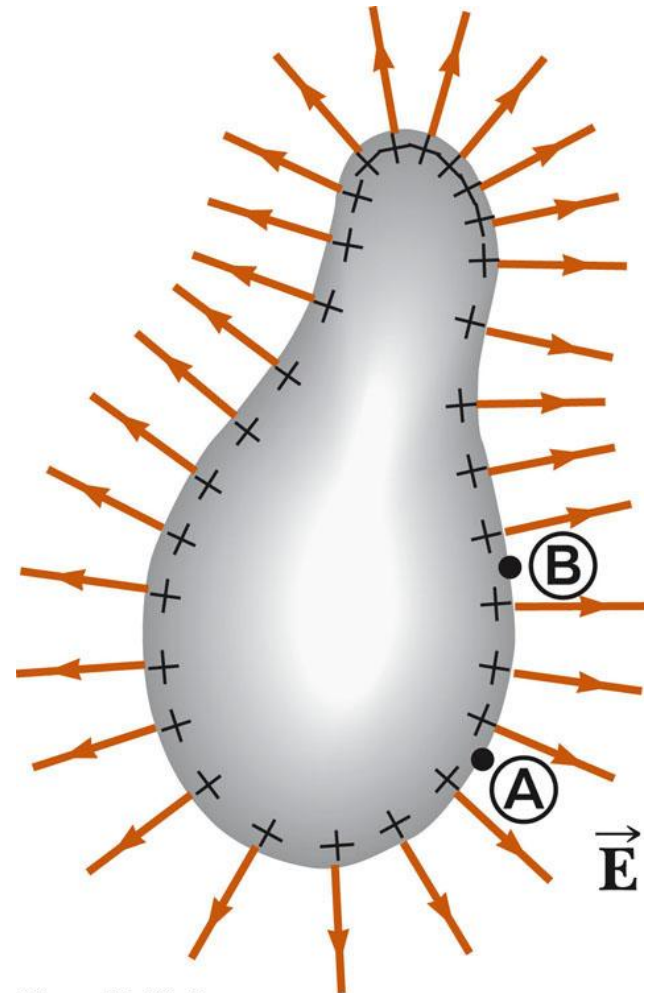
$$V = \frac{k_e Q}{\ell} \ln \left(\frac{\ell + \sqrt{a^2 + \ell^2}}{a} \right)$$



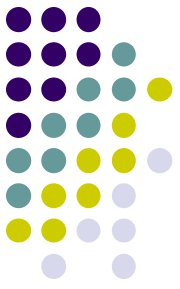


V Due to a Charged Conductor

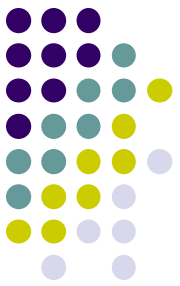
- Consider two points on the surface of the charged conductor as shown
- \vec{E} is always perpendicular to the displacement $d\vec{s}$
- Therefore, $\vec{E} \cdot d\vec{s} = 0$
- Therefore, the potential difference between A and B is also zero



V Due to a Charged Conductor, cont.

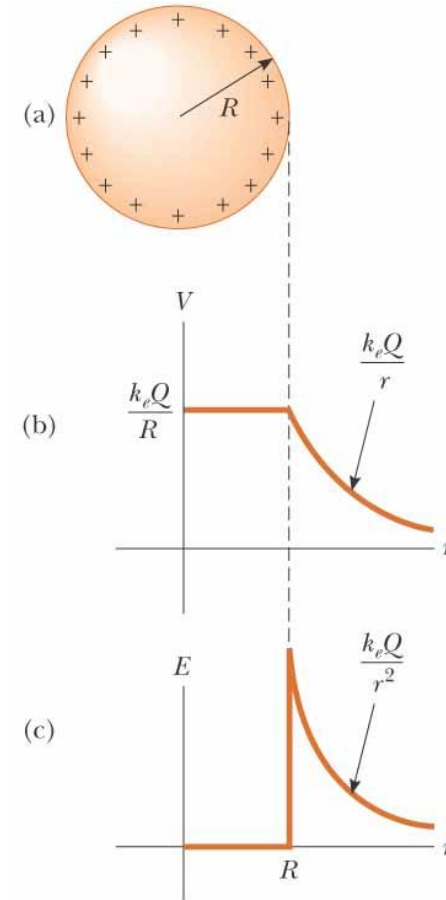


- V is constant everywhere on the surface of a charged conductor in equilibrium
 - $\Delta V = 0$ between any two points on the surface
- The surface of any charged conductor in electrostatic equilibrium is an equipotential surface
- Because the electric field is zero inside the conductor, we conclude that the electric potential is constant everywhere inside the conductor and equal to the value at the surface



E Compared to V

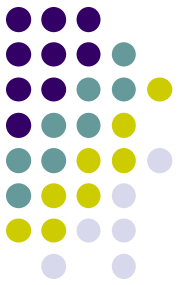
- The electric potential is a function of r
- The electric field is a function of r^2
- The effect of a charge on the space surrounding it:
 - The charge sets up a vector electric field which is related to the force
 - The charge sets up a scalar potential which is related to the energy



Irregularly Shaped Objects

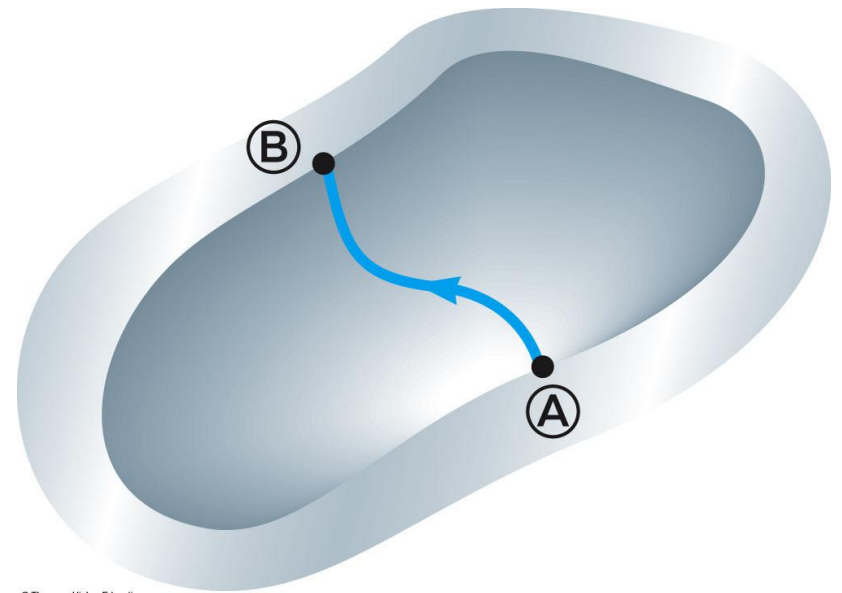


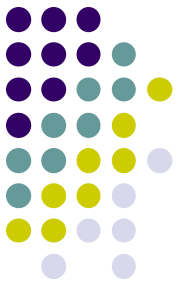
- The charge density is high where the radius of curvature is small
 - And low where the radius of curvature is large
- The electric field is large near the convex points having small radii of curvature and reaches very high values at sharp points



Cavity in a Conductor

- Assume an irregularly shaped cavity is inside a conductor
- Assume no charges are inside the cavity
- The electric field inside the conductor must be zero





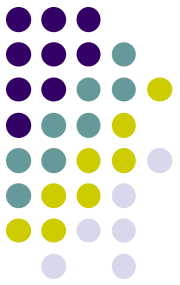
Cavity in a Conductor, cont

- The electric field inside does not depend on the charge distribution on the outside surface of the conductor

- For all paths between A and B ,

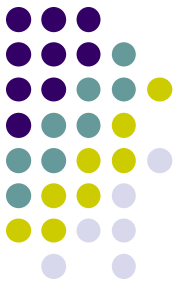
$$V_B - V_A = -\int_A^B \vec{\mathbf{E}} \cdot d\vec{\mathbf{s}} = 0$$

- A cavity surrounded by conducting walls is a field-free region as long as no charges are inside the cavity



Corona Discharge

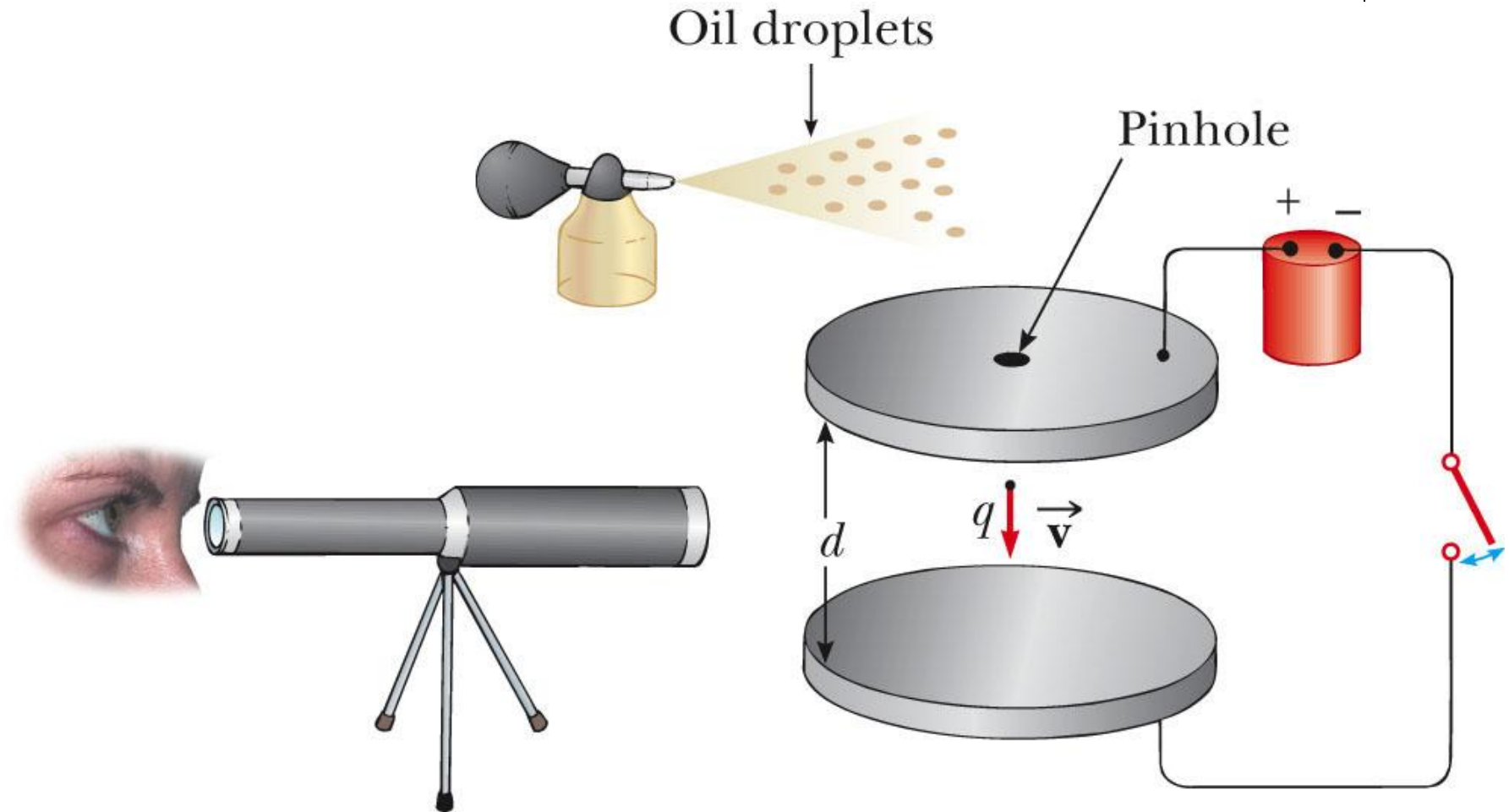
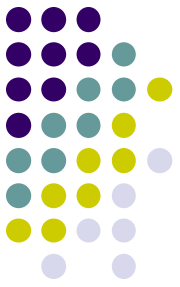
- If the electric field near a conductor is sufficiently strong, electrons resulting from random ionizations of air molecules near the conductor accelerate away from their parent molecules
- These electrons can ionize additional molecules near the conductor



Corona Discharge, cont.

- This creates more free electrons
- The **corona discharge** is the glow that results from the recombination of these free electrons with the ionized air molecules
- The ionization and corona discharge are most likely to occur near very sharp points

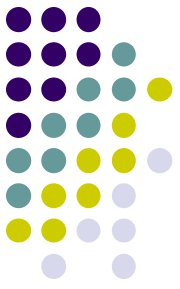
Millikan Oil-Drop Experiment – Experimental Set-Up



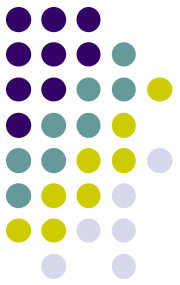
Telescope with
scale in eyepiece

**PLAY
ACTIVE FIGURE**

Millikan Oil-Drop Experiment



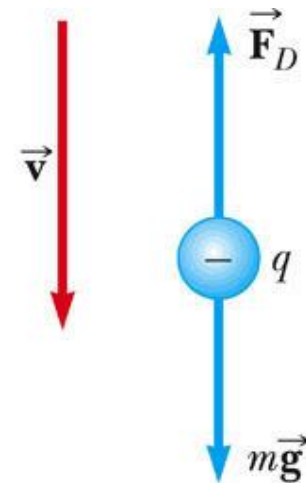
- Robert Millikan measured e , the magnitude of the elementary charge on the electron
- He also demonstrated the quantized nature of this charge
- Oil droplets pass through a small hole and are illuminated by a light



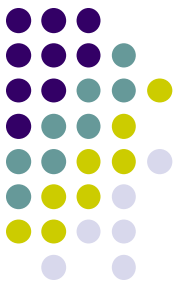
Oil-Drop Experiment, 2

- With no electric field between the plates, the gravitational force and the drag force (viscous) act on the electron
- The drop reaches terminal velocity with

$$\vec{F}_D = m\vec{g}$$

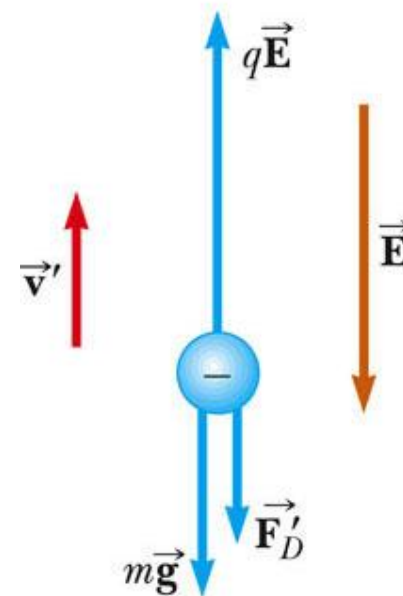


(a)



Oil-Drop Experiment, 3

- When an electric field is set up between the plates
 - The upper plate has a higher potential
- The drop reaches a new terminal velocity when the electrical force equals the sum of the drag force and gravity



(b)

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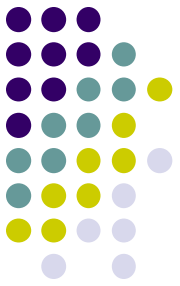
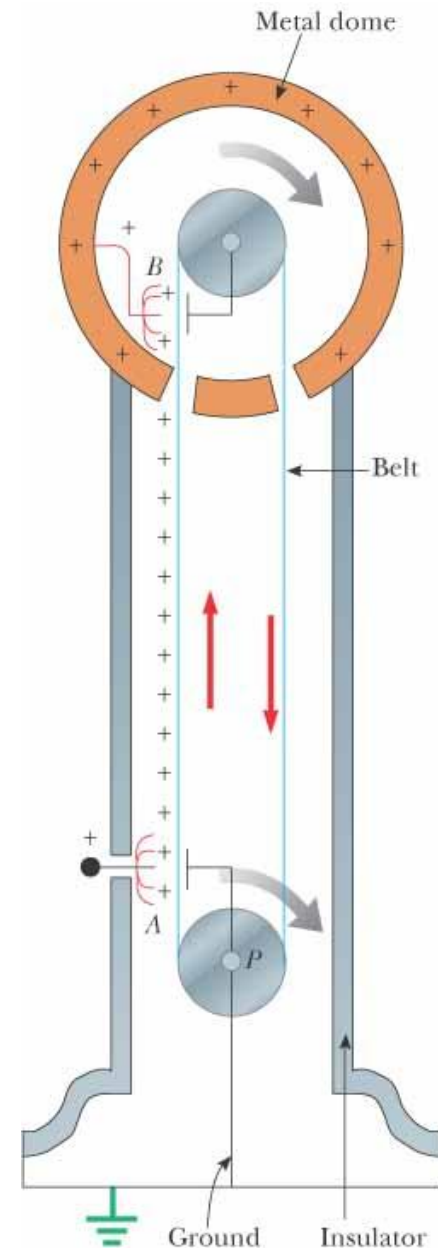


Oil-Drop Experiment, final

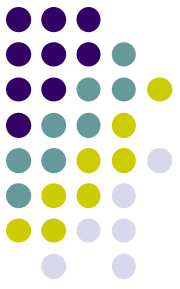
- The drop can be raised and allowed to fall numerous times by turning the electric field on and off
- After many experiments, Millikan determined:
 - $q = ne$ where $n = 0, -1, -2, -3, \dots$
 - $e = 1.60 \times 10^{-19} \text{ C}$
- This yields conclusive evidence that charge is quantized
- Use the active figure to conduct a version of the experiment

Van de Graaff Generator

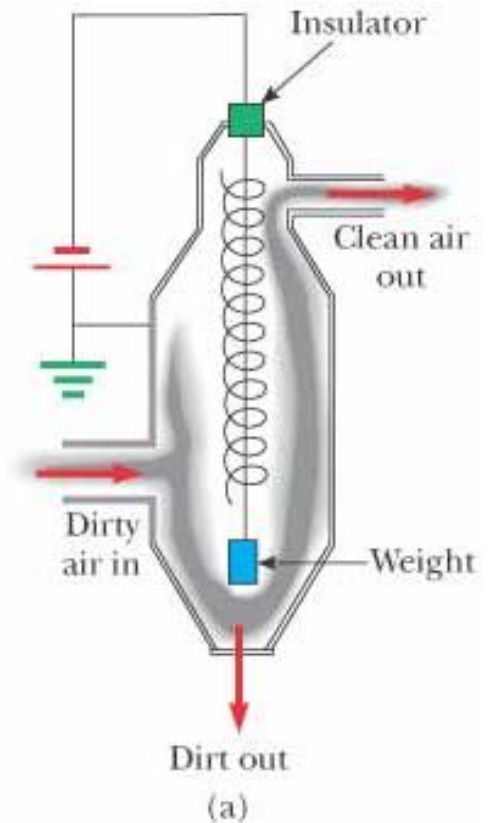
- Charge is delivered continuously to a high-potential electrode by means of a moving belt of insulating material
- The high-voltage electrode is a hollow metal dome mounted on an insulated column
- Large potentials can be developed by repeated trips of the belt
- Protons accelerated through such large potentials receive enough energy to initiate nuclear reactions



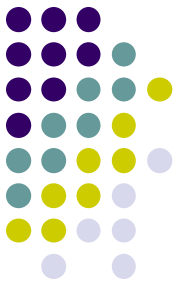
Electrostatic Precipitator



- An application of electrical discharge in gases is the electrostatic precipitator
- It removes particulate matter from combustible gases
- The air to be cleaned enters the duct and moves near the wire
- As the electrons and negative ions created by the discharge are accelerated toward the outer wall by the electric field, the dirt particles become charged
- Most of the dirt particles are negatively charged and are drawn to the walls by the electric field

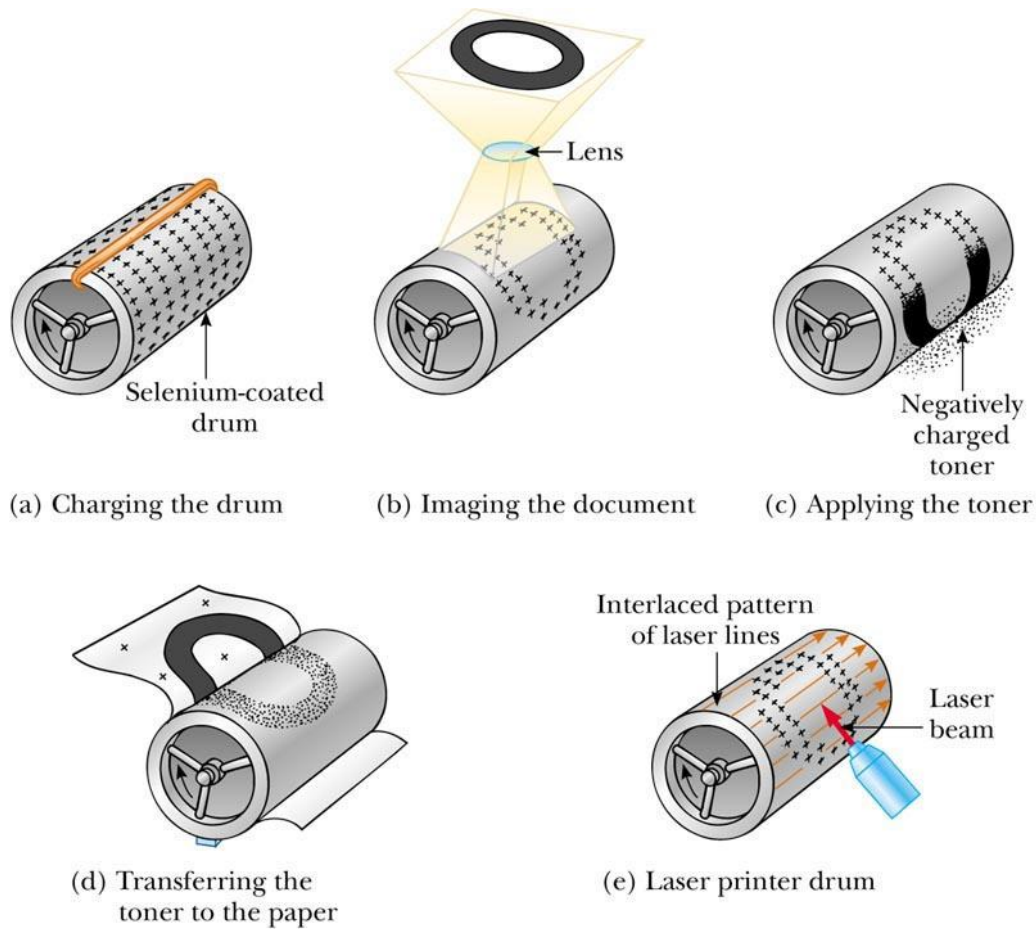


Application – Xerographic Copiers



- The process of xerography is used for making photocopies
- Uses photoconductive materials
 - A photoconductive material is a poor conductor of electricity in the dark but becomes a good electric conductor when exposed to light

The Xerographic Process



Application – Laser Printer



- The steps for producing a document on a laser printer is similar to the steps in the xerographic process
- A computer-directed laser beam is used to illuminate the photoconductor instead of a lens